

Remarks on QPE:

$|\psi\rangle$ eigenvector of U unitary, $U|\psi\rangle = e^{2\pi i\theta}|\psi\rangle$, remember the QPE algorithm.

Question: what if $|\psi\rangle$ is unknown?

Ex.: if $|\psi\rangle = \alpha_0|\psi_0\rangle + \alpha_1|\psi_1\rangle$ with $|\alpha_0|^2 + |\alpha_1|^2 = 1$ and $U|\psi_0\rangle = e^{2\pi i\theta_0}|\psi_0\rangle$, $U|\psi_1\rangle = e^{2\pi i\theta_1}|\psi_1\rangle$, $\langle\psi_0|\psi_1\rangle = 0$.

When we apply the algorithm we get $\frac{1}{2^t} \sum_{l=0}^{2^t-1} \sum_{k=0}^{2^t-1} \left(e^{2\pi i(\theta_0 - \frac{l}{2^t})k} \alpha_0 |l\rangle \otimes |\psi_0\rangle + e^{2\pi i(\theta_1 - \frac{l}{2^t})k} \alpha_1 |l\rangle \otimes |\psi_1\rangle \right)$.

When we measure I/O register, we obtain $l \in \{0, 1, \dots, 2^t-1\}$ with probability $P(l) = \left\| \frac{1}{2^t} \sum_{k=0}^{2^t-1} \left(e^{2\pi i(\theta_0 - \frac{l}{2^t})k} \alpha_0 |\psi_0\rangle + e^{2\pi i(\theta_1 - \frac{l}{2^t})k} \alpha_1 |\psi_1\rangle \right) \right\|^2$
 $= \frac{1}{2^{2t}} |\alpha_0|^2 \left| \sum_{k=0}^{2^t-1} e^{2\pi i(\theta_0 - \frac{l}{2^t})k} \right|^2 + \frac{1}{2^{2t}} |\alpha_1|^2 \left| \sum_{k=0}^{2^t-1} e^{2\pi i(\theta_1 - \frac{l}{2^t})k} \right|^2$.

This situation arises in Quantum Counting (when we want to find m in Grover's algorithm).

Harrow-Hassidim-Lloyd (HHL) algorithm (2008)

Solving linear systems of equations using quantum computers

Problem: given $A \in \mathbb{C}^{N \times N}$, $b \in \mathbb{C}^N$, find $x \in \mathbb{C}^N$ s.t. $Ax = b$.

Gaussian elimination: $O(N^3)$.

Recall the conditioning number of A is $\kappa(A) = \|A\| \cdot \|A^{-1}\|$; if

A is positive and hermitian, $\kappa(A) = \frac{\lambda_{\max}}{\lambda_{\min}}$.

A is s -sparse if every row of A contains at most s elements $\neq 0$. To solve the problem in this case we can use the conjugate gradient algorithm with $O(N s \kappa(A) \log(1/\epsilon))$.

The HHL algorithm, on A positive s -sparse, takes only $O((\log N) s^2 \kappa(A)^2 / \epsilon)$ quantum steps.

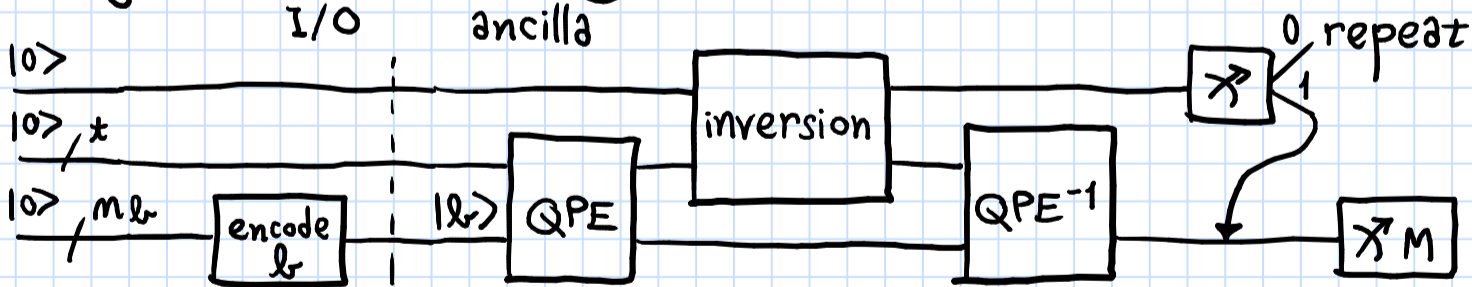
Actually we get a state $\frac{|x\rangle}{\|x\|}$ so we can measure for any observable $M \in \mathbb{C}^{N \times N}$ the quantity $(M)_{\frac{|x\rangle}{\|x\|}} = \frac{\langle x | M | x \rangle}{\|x\|^2}$.

If A is not s -sparse (Vossing, 2018): $O(\sqrt{N} \log N \kappa(A)^2)$.

Two pillars of the HHL algorithm:

- 1) we use QPE on $U = e^{2\pi i A}$, so we need an efficient implementation of controlled U (hamiltonian simulation);
- 2) there is a "branch selection" step, we measure the ancilla and according to the outcome (1 or 0) we perform further quantum operations on I/O register.

Register $\underbrace{|H^{m_b}\rangle}_{\text{I/O}} \otimes \underbrace{|H^t\rangle}_{\text{ancilla}} \otimes |H\rangle \rightarrow$ we measure this



Inversion gate acts on $|H^t\rangle \otimes |H\rangle$ on the basis:

$$|l\rangle \otimes |0\rangle \mapsto |l\rangle \otimes (f(l)|1\rangle + \sqrt{1-f(l)^2}|0\rangle),$$

$$|l\rangle \otimes |1\rangle \mapsto |l\rangle \otimes (-\sqrt{1-f(l)^2}|1\rangle + f(l)|0\rangle),$$

$$\text{where } f(l) = \begin{cases} \frac{\lambda_{\min}}{l/2^t} = \frac{c}{l} & \text{if } l \geq c \\ 0 & \text{otherwise} \end{cases}$$

We assume $N = 2^{m_b}$ and we have to normalize b ($\|b\|=1$).

Now we do the calculations of the algorithm. To simplify, we do the case where all eigenvalues of A are of the form $\frac{l}{2^t}$.

When we measure the ancilla qubit, $P(1) = \sum_{k=0}^{N-1} |b_k|^2 \left(\frac{\lambda_{\min}}{\lambda_k} \right)^2 \frac{1}{2^t}$ where λ_k are the eigenvalues of A and b_k are the coefficients of b in the base of eigenstates of A , so $P(1) \geq \|b\|^2 (\kappa(A))^{-2}$.

After measuring, I/O register is in state $|x\rangle / \|x\|$.